

GRB Prompt Emission Spectra: The Synchrotron Revenge

Maria Edvige Ravasio^{1 2}

¹ Università degli Studi di Milano-Bicocca, Dipartimento di Fisica U2, Piazza della Scienza, 3, I20126, Milano, Italy

² INAF Osservatorio Astronomico di Brera, via Bianchi 46, I23807 Merate (LC), Italy

E-mail: m.ravasio5@campus.unimib.it, maria.ravasio@inaf.it

ABSTRACT

After more than 40 years from their discovery, the long-lasting tension between predictions and observations of GRBs prompt emission spectra starts to be solved. We found that the observed spectra can be produced by the synchrotron process, if the emitting particles do not completely cool. Evidence for incomplete cooling was recently found in Swift GRBs spectra with prompt observations down to 0.5 keV (Oganesyan et al. 2017, 2018), characterized by an additional low-energy break. In order to search for this break at higher energies, we analysed the 10 long and 10 short brightest GRBs detected by the Fermi satellite in over 10 years of activity. We found that in 8/10 long GRBs there is compelling evidence of a low energy break (below the peak energy) and the photon indices below and above that break are remarkably consistent with the values predicted by the synchrotron spectrum ($-2/3$ and $-3/2$, respectively). None of the ten short GRBs analysed shows a break, but the low energy spectral slope is consistent with $-2/3$. Within the framework of the GRB standard model, these results imply a very low magnetic field in the emission region, at odds with expectations. I also present the spectral evolution of GRB 190114C, the first GRB detected with high significance by the MAGIC Telescopes, which shows the compresence (in the keV-MeV energy range) of the prompt and of the afterglow emission, the latter rising and dominating the high energy part of the spectral energy range.

KEY WORDS: workshop: proceedings — prompt emission – gamma-ray burst – radiation mechanism

1. Introduction

The nature of the radiation mechanism behind the prompt emission of Gamma-Ray Bursts (GRBs) has puzzled astronomers since their discovery and still remains not fully understood after more than 40 years of observations. The non-thermal shape of the observed spectra and the likely presence of accelerated particles in a magnetized region lead to the suggestion that the synchrotron process could be the responsible for such γ -ray emission (Rees & Meszaros 1994; Katz 1994; Tavani 1996; Sari & Piran 1997). However, the inconsistency between the observed GRB spectra and the theoretical spectral shapes expected for synchrotron emission has kept the debate alive for many years.

From an observational point of view, the typical GRB prompt emission spectrum is usually fitted with the so-called Band function (Band et al. 1993), namely two power-laws with slopes α and β , smoothly connected at the peak energy E_{peak} of the $\nu F\nu$ spectrum. The slope α of the low-energy power-law has been found to have an average value of $\langle\alpha\rangle \sim -1$ for long bursts, while short GRBs, on average, are harder: $\langle\alpha\rangle \sim -0.4$ (Preece et al. 2000; Kaneko et al. 2006; Ghirlanda et al. 2009; Nava et

al. 2011; Goldstein et al. 2012; Gruber et al. 2014).

For the typical physical parameters expected in the prompt GRB emitting region, the electrons should radiate in a regime of fast cooling (Ghisellini et al. 2000) and the observed spectrum should have a photon slope $\alpha = -3/2$ below the peak energy E_{peak} . Thus, most of the typical observed slopes are inconsistent with the synchrotron predictions, since the majority of GRBs are harder than $-3/2$, and in some case even harder than $-2/3$, corresponding to the limiting case of the single particle spectrum, the so-called synchrotron 'line-of-death' (Preece et al. 1998). The observed hard value of the low-energy photon index is one of the key observational features that strongly challenged the synchrotron interpretation.

Recently, extending the investigation of the prompt emission of 34 long GRBs down to the soft X-rays, Oganesyan et al. (2017) found low-energy breaks in their spectra. The distribution of these breaks is centered around $E_{break} \sim 30$ keV and the slopes below and above that break are distributed around $\langle\alpha_1\rangle = -0.51$ ($\sigma = 0.24$) and $\langle\alpha_2\rangle = -1.56$ ($\sigma = 0.26$). These slopes are consistent, within 1σ , with the expectations of the synchrotron the-

ory ($\alpha_1 = -2/3$ and $\alpha_2 = -3/2$) in a marginally fast cooling regime (i.e., when $\nu_{cool} \sim \nu_{min}$ (Daigne et al. 2011; Beniamini & Piran 2013)). Since this low-energy break was found in the soft X-rays, we decided to extend the search of this feature also at higher energies through GRBs detected by the Fermi/GBM instrument (8 keV - 40 MeV).

We analyzed GRB 160625B (Ravasio et al. 2018), a very bright burst whose light curve is composed by three distinct emission episodes: a precursor, the main event and a last dimmer event. We performed a time-resolved spectral analysis on the main event using an empirical function composed of three power-laws smoothly connected at two breaks. Fitting with three power-laws actually provides an improvement of the fit (with respect to a function with only two power-laws) in the majority of the spectra analysed ($\sigma \geq 8$, according to the F-test). The time-resolved spectra of this burst are characterised by a low-energy power-law photon index centered around $\alpha_1 = -0.63$ ($\sigma=0.08$), the presence of an additional low-energy break at $E_{break} \sim 50$ -100 keV, a second power-law photon index centered around $\alpha_2 = -1.48$ ($\sigma=0.09$) between the break and the peak energy, a second spectral break (representing the peak in νF_ν) varying with time in the range $E_{peak} \sim 300$ keV - 6 MeV, and a third power-law photon index $\beta \sim -2.6$ describing the spectrum above the peak energy. The slopes below and above the break are remarkably consistent with the synchrotron predicted values.

Motivated by these results, we performed a systematic search for this feature in a larger sample of GRBs (Ravasio et al. 2019). To this aim, we selected the 10 brightest long GRBs and the 10 brightest short GRBs detected by the GBM instrument over 10 years of activity.

For long GRBs, the time-integrated analysis shows that in 8/10 of the brightest GBM bursts, the standard fitting function fails to provide an acceptable fit: data require an additional spectral break E_{break} , located between ~ 10 keV and 300 keV. For these eight GRBs, a detailed time-resolved analysis revealed that $\sim 70\%$ of the spectra analysed also show strong evidence of an additional low-energy spectral break. The photon index of the power law below E_{break} is distributed around $\langle\alpha_1\rangle = -0.58$ ($\sigma=0.16$) while the power law between E_{break} and E_{peak} has photon index $\langle\alpha_2\rangle = -1.52$ ($\sigma=0.20$), both remarkably consistent with the predicted values for synchrotron emission (see Fig. 1 for an example of a typical spectrum).

The remaining time-resolved spectra ($\sim 30\%$) are best fitted by the standard function. In these cases, one power law is sufficient to model the spectra below E_{peak} , and the mean value of its photon index is $\langle\alpha\rangle = -1.02$ ($\sigma=0.19$). This value is in agreement with previous works and, interestingly, it lies between the values of α_1 and α_2 , as shown in Fig. 2 by the black empty his-

togram. Speculating that most of the spectra show a break below the peak energy, the value of α can be understood as a sort of average value between α_1 and α_2 . The consistency with synchrotron emission in the long GRBs spectra is supported also by the recent works of Oganessian et al. (2019) and Burgess et al. (2019), where a physical synchrotron model has been found to fit well the spectra. Note, however, that the majority of spectra analysed by Burgess et al. (2019) displays a regime of slow cooling, implying a even lower efficiency of radiation.

Regarding short GRBs, our analysis shows that no spectra shows a second break at low energies. They are best-fitted by the standard fitting function, namely with only one component below the peak energy. The low-energy slope of this component is centered around $\langle\alpha\rangle = -0.78$ ($\sigma=0.23$). As for α_1 in long GRBs, this value is consistent with the low-energy synchrotron photon index $-2/3$.

These results suggest that the underlying population of electrons does not cool completely. In fact, let us assume a power law distribution of electrons, injected above some injection energy γ_{min} : they cool via synchrotron process (and the observed slope $-3/2$ below E_{peak} is the footprint of their cooling), but then at some energy γ_{cool} they stop cooling, and the single electron spectral slope $-2/3$ appears. Therefore, interpreting E_{break} as the synchrotron cooling frequency ν_{cool} , the implied magnetic field B' in the (comoving) emitting region is small, i.e. between 1 and 40 G. The situation is even more extreme in short GRBs. Contrary to long GRBs, their spectra are described by only one power law below E_{peak} , with slope $\sim -2/3$, suggesting that $\nu_{cool} \simeq \nu_{min}$ and the cooling is even more incomplete, thus yielding an even smaller estimate of B' .

In the framework of the GRB standard model, the expected magnetic field, in the emitting region, should be of the order of 10^6 G, i.e. much larger than the estimates above. If the emitting region were at a standard distance ($R \sim 10^{13}$ cm), a small B' would imply a strong inverse Compton component, whose luminosity would exceed the synchrotron one by a factor of $\sim 10^7$. A weak magnetic field could be possible if the emitting region were at larger distances (i.e. 10^{18} cm) where, however, the afterglow emission could also take place, and the expected variability timescale would be much larger than the typical short values observed in the prompt emission.

Recently, Ghisellini et al. (2019) proposed that spectral results of Oganessian et al. (2017), Oganessian et al. (2018), Ravasio et al. (2018) and Ravasio et al. (2019) could be explained considering synchrotron emission from protons rather than electrons. Due to their larger mass, the synchrotron cooling timescale of protons can be relatively long, i.e. $\sim 10^8$ times longer than

electrons. This can account for $\nu_{cool} \sim 100$ keV with a standard magnetic field of $B' \sim 10^6$ G and considering the emission region at $R \sim 10^{13}$ cm (thus accounting for a short timescale variability).

2. The γ -ray extra-component in GRB 190114C

GRB 190114C is the first GRB detected with high significance at TeV energies by the MAGIC telescopes (Mirzoyan et al. 2019). It was also detected by Fermi/GBM with a [10-1000 keV] fluence $f = 4.433 \pm 0.005 \times 10^{-4}$ erg/cm². We performed (Ravasio et al. 2019) a detailed time-resolved spectral analysis of the Fermi/GBM emission, from 10 keV to 40 MeV, up to ~ 60 s from the trigger time.

We found that the first 4.8 seconds are characterized by the standard prompt spectral shape, fitted by a standard function with typical parameters. After 4.8 seconds from the trigger, there is evidence of the emergence of another spectral component, superimposed on the typical prompt shape spectrum. This component rises and decays quickly, reaching its peak flux ($F = 1.7 \times 10^{-5}$ erg/cm²/s, in the 10 keV-40 MeV energy range) at ~ 6 s. We found that this additional component is well fitted by a power law of spectral index $\Gamma_{PL} \sim -2$ (see Fig. 1 in Ravasio et al. (2019)).

We interpreted this component as the afterglow of the burst, because i) it appears after the trigger of the prompt event, and peaks when most of the prompt emission energy has already been radiated, ii) it lasts much longer than the prompt emission, iii) it is characterized by a spectral index ($\Gamma_{PL} \sim -2$) typical of known afterglows, iv) with the exception of the early variable phases, its light curve smoothly decays as $F \propto t^{-1}$, typical of the known afterglows.

Therefore, we found the evidence of the co-presence in the keV-MeV energy range of both the prompt and the afterglow components, already after ~ 5 s from trigger time. We interpret the peak of the afterglow as the fireball deceleration time and estimate the bulk Lorentz factor during the coasting phase $\Gamma_0=700$ or 130, (assuming an homogeneous or wind medium, respectively).

References

Band, D., Matteson, J., Ford, L., et al. 1993, ApJ, 413, 281
 Beniamini, P., & Piran, T. 2013, ApJ, 769, 69
 Burgess, J. M., Bégúé, D., Greiner, J., et al. 2019, Nature Astronomy, 471
 Daigne, F., Bošnjak, Ž., & Dubus, G. 2011, A&A, 526, A110
 Ghirlanda, G., Nava, L., Ghisellini, G., et al. 2009, A&A, 496, 585
 Ghisellini, G., Celotti, A., & Lazzati, D. 2000, MNRAS, 313, L1

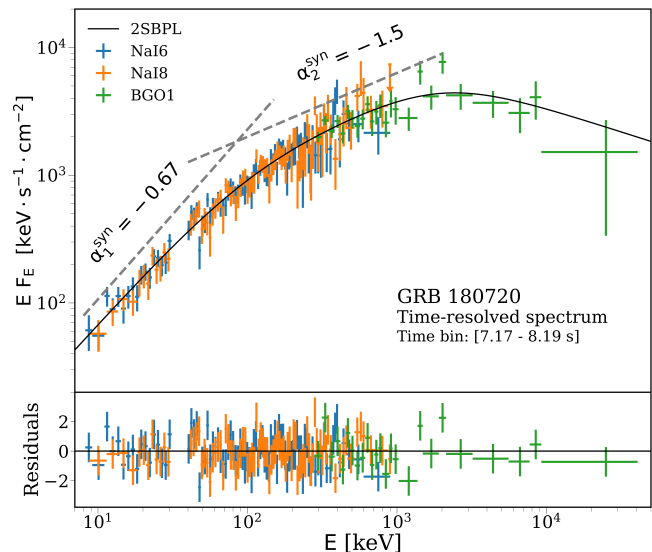


Fig. 1. Example of a spectrum best fitted by a 2SBPL (i.e. three power laws smoothly connected at two breaks). The best fit values of the 2SBPL model parameters $\alpha_1 = -0.71$, $E_{break} = 93.6$ keV, $\alpha_2 = -1.47$, $E_{peak} = 2.4$ MeV, and $\beta = -2.38$. The two dashed lines show the power laws with the photon indices predicted by synchrotron emission. Data-to-model residuals are shown in the bottom panel. From Ravasio et al. (2019).

Ghisellini, G., Ghirlanda, G., Oganessian, G., et al. 2019, arXiv e-prints, arXiv:1912.02185
 Goldstein, A., Burgess, J. M., Preece, R. D., et al. 2012, ApJS, 199, 19
 Gruber, D., Goldstein, A., Weller von Ahlefeld, V., et al. 2014, ApJS, 211, 12
 Kaneko, Y., Preece, R. D., Briggs, M. S., et al. 2006, ApJS, 166, 298
 Katz, J. I. 1994, ApJL, 432, L107
 Mirzoyan, R., Noda, K., Moretti, E., et al. 2019, GRB Coordinates Network 23701, 1
 Nava, L., Ghirlanda, G., Ghisellini, G., et al. 2011, A&A, 530, A21
 Oganessian, G., Nava, L., Ghirlanda, G., et al. 2017, ApJ, 846, 137
 Oganessian, G., Nava, L., Ghirlanda, G., et al. 2018, A&A, 616, A138
 Oganessian, G., Nava, L., Ghirlanda, G., et al. 2019, A&A, 628, A59
 Preece, R. D., Briggs, M. S., Mallozzi, R. S., et al. 1998, ApJL, 506, L23
 Preece, R. D., Briggs, M. S., Mallozzi, R. S., et al. 2000, ApJS, 126, 19
 Ravasio, M. E., Oganessian, G., Ghirlanda, G., et al. 2018, A&A, 613, A16
 Ravasio, M. E., Ghirlanda, G., Nava, L., et al. 2019, A&A, 625, A60

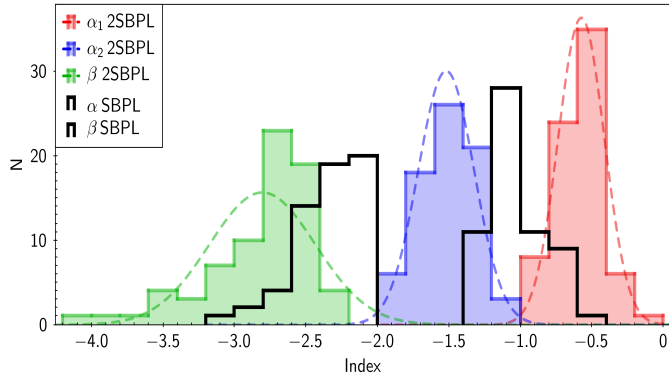


Fig. 2. Distribution of the spectral indices for the time-resolved fits of the 8 long GRBs showing an additional spectral break. The spectral indices α_1 , α_2 and β of the 2SBPL model are shown with red, blue, and green filled histograms, respectively. Gaussian functions showing the central value and standard deviation of the distributions are overlapped to the histograms (colour-coded dashed curves). The black empty histograms represent the distributions of the two photon indices α and β of the SBPL model, for spectra where the SBPL is the best fit model. From Ravasio et al. (2019).

Ravasio, M. E., Oganessian, G., Salafia, O. S., et al. 2019, *A&A*, 626, A12
 Rees, M. J., & Meszaros, P. 1994, *ApJL*, 430, L93
 Sari, R., & Piran, T. 1997, *MNRAS*, 287, 110
 Tavani, M. 1996, *ApJ*, 466, 768