

# Fixed point indices of central configurations

D.L. Ferrario

December 19, 2014

## Abstract

Central configurations of  $n$  point particles in  $E \cong \mathbb{R}^d$  with respect to a potential function  $U$  are shown to be the same as the fixed points of the normalized gradient map  $F = -\nabla_M U / \|\nabla_M U\|_M$ , which is an  $SO(d)$ -equivariant self-map defined on the inertia ellipsoid. We show that the  $SO(d)$ -orbits of fixed points of  $F$  are all fixed points of the map induced on the quotient by  $SO(d)$ , and give a formula relating their indices (as fixed points) with their Morse indices (as critical points). At the end, we give an example of a non-planar relative equilibrium which is not a rotating central configuration.

*MSC Subject Class:* 55M20, 37C25, 70F10.

*Keywords:* Central configurations, relative equilibria,  $n$ -body problem.

## 1 Central configurations

Let  $E = \mathbb{R}^d$  be the  $d$ -dimensional Euclidean space, and  $X = E^n \setminus \Delta$  the *configuration space* of  $E$ , defined as the space of all  $n$ -tuples of distinct points  $\mathbf{q} = (\mathbf{q}_1, \dots, \mathbf{q}_n) \in E^n$  such that  $\mathbf{q} \notin \Delta$ , where  $\Delta$  is the collision set

$$\Delta = \bigcup_{i < j} \{\mathbf{q} \in E^n : \mathbf{q}_i = \mathbf{q}_j\}.$$

Let  $U: X \rightarrow \mathbb{R}$  be a regular potential function. For example, the gravitational potential

$$(1.1) \quad U(\mathbf{q}) = \sum_{i < j} \frac{m_i m_j}{\|\mathbf{q}_i - \mathbf{q}_j\|^\alpha},$$

or more general potentials (charged particles, ...).

Given  $n$  real non-zero numbers  $m_1, \dots, m_n$  (representing the *masses* of the  $n$  interacting point particles), let  $\langle -, - \rangle_M$  denote the *mass-metric* on  $E^n$ , defined as  $\langle \mathbf{v}, \mathbf{w} \rangle_M = \sum_{j=1}^n m_j \mathbf{v}_j \cdot \mathbf{w}_j$ , where  $\mathbf{v}, \mathbf{w} \in \mathbb{R}^{dn}$  are two vectors tangent to  $E^n$ , and  $\cdot$  denotes the standard scalar product in  $E$ . If the masses are positive, then the mass-metric is a non-degenerate scalar product

on  $E^n$ , which yields both the kinetic quadratic form (on the tangent bundle)  $2K = \|\dot{\mathbf{q}}\|_M^2 = \sum_i m_i \|\dot{\mathbf{q}}_i\|^2$  and the inertia form  $I = \|\mathbf{q}\|_M^2 = \sum_i m_i \|\mathbf{q}_i\|^2$ .

Given such  $m_i$ , let  $\nabla_M$  denote the gradient of  $U$  with respect to the bilinear product  $\langle -, - \rangle_M$ . Let us recall that if  $dU(\mathbf{q})$  denotes the differential of  $U$  evaluated in  $\mathbf{q}$ , then for each tangent vector  $\mathbf{v} \in T_{\mathbf{q}}E^n$ ,  $dU(\mathbf{q})[\mathbf{v}] = \langle \nabla_M U(\mathbf{q}), \mathbf{v} \rangle_M$ , and therefore

$$dU(\mathbf{q})[\mathbf{v}] = \sum_j \frac{\partial U}{\partial \mathbf{q}_j} \cdot \mathbf{v}_j = \sum_j m_j \left( m_j^{-1} \frac{\partial U}{\partial \mathbf{q}_j} \right) \cdot \mathbf{v}_j,$$

from which it follows that in standard coordinates  $(\nabla_M U)_j = m_j^{-1} \frac{\partial U}{\partial \mathbf{q}_j}$ .

Given the mass-metric gradient  $\nabla_M U$ , the corresponding *Newton equations* are

$$(1.2) \quad \frac{d^2 \mathbf{q}}{dt^2} = \nabla_M U(\mathbf{q}).$$

**(1.3) Definition.** A configuration  $\mathbf{q} \in X$  is a *central configuration* iff there exists  $\lambda \in \mathbb{R}$  such that  $\nabla_M U(\mathbf{q}) = \lambda \mathbf{q}$ .

(1.4) *Remark.* If  $U$  is homogeneous, this is equivalent to:  $\mathbf{q}$  is a central configuration iff there is a real-valued function  $\phi(t)$  such that  $\phi(t)\mathbf{q}$  is a solution of (1.2)<sup>1</sup>. If  $U$  is homogeneous, the set of central configurations is a cone in  $X$ .

(1.5) *Remark.* Furthermore, if  $U$  is invariant with respect to the group of all translations of  $E$ , then central configurations belong to the subspace

$$Y = \left\{ \mathbf{q} \in X : \sum_j m_j \mathbf{q}_j = 0 \right\} \subset X \subset E^n,$$

and  $\forall \mathbf{q} \in X \implies \nabla_M U(\mathbf{q}) \in \bar{Y} \subset E^n$ , where  $\bar{Y}$  is the closure of  $Y$  in  $E^n$ . Therefore, if  $U$  is translation-invariant, the set of central configurations is a subset of  $Y$ . Sometimes central configurations are defined with the equation  $\nabla_M U(\mathbf{q}) = \lambda(\mathbf{q} - \mathbf{c})$ , where  $\mathbf{c}$  is the center of mass  $\mathbf{c} = (\sum_j m_j)^{-1} \sum_j m_j \mathbf{q}_j$  of the configuration  $\mathbf{q}$ . This equation is invariant with respect to translations (if  $U$  is so).

**(1.6) Definition.** A configuration  $\mathbf{q} \in X$  is a *relative equilibrium* iff there is a one-parameter group of rotations  $\varphi^t: E \rightarrow E$  (around the origin, without loss of generality) such that

$$\varphi^t(\mathbf{q}) = (\varphi^t(\mathbf{q}_1), \varphi^t(\mathbf{q}_2), \dots, \varphi^t(\mathbf{q}_n))$$

satisfies the equations of motion (1.2).<sup>2</sup>

---

<sup>1</sup>Cf. [8], p. 61.

<sup>2</sup>Cf. [17], p. 47.

One-parameter subgroups of  $SO(N)$  are of the form  $\varphi^t(\mathbf{q}_1) = e^{t\Omega}\mathbf{q}_1$ , with  $\Omega$  skew-symmetric  $N \times N$  (non-zero) matrix.<sup>3</sup> If  $U$  is invariant with respect to the above-mentioned one-parameter group of rotations  $\varphi^t = e^{t\Omega}$ , a relative equilibrium satisfies the equation

$$\Omega^2\mathbf{q} = \nabla_M U(\mathbf{q}) .$$

It follows that if  $\dim(E) = 2$ , then  $\Omega = \begin{bmatrix} 0 & -\omega \\ \omega & 0 \end{bmatrix}$  with  $\omega \in \mathbb{R}$ ,  $\omega \neq 0$ , and therefore such a relative equilibrium configuration is a central configuration,  $-\omega^2\mathbf{q} = \nabla_M U(\mathbf{q})$ . Conversely, planar central configurations with  $U(\mathbf{q}) > 0$  (if  $U$  is homogeneous of negative degree) yield relative equilibria, with a suitable (angular speed)  $\omega$ .

If  $\dim(E) = 3$ , then since the non-zero  $3 \times 3$  skew-symmetric matrix  $\Omega$  has rank 2,  $E$  can be written as  $\ker \Omega \oplus E'$ , where  $\ker \Omega$  is the fixed direction of the rotations  $\varphi^t = e^{t\Omega}$ , and  $E'$  is the orthogonal complement of  $\ker \Omega$ . In a suitable reference,  $\Omega = \begin{bmatrix} 0 & -\omega & 0 \\ \omega & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$ . For further reference, let

$$(1.7) \quad P: E \rightarrow E'$$

denote the orthogonal projection.

If  $U$  is the homogeneous Newtonian potential of (1.1) with  $\alpha > 0$  and  $m_j > 0$ , then it is easy to see that relative equilibrium configurations must belong to the plane  $E'$ . This is not true in general: it is possible to find examples of relative equilibria which are not planar – see (4.1) (and hence they are not central configurations). For more on equilibrium (and homographic) solutions: [19] (§369–§382bis at pp. 284–306), [4], [2].

Recent and non recent relevant literature on central configurations: [14], [7], [16, 17], [10, 11, 12], [8], [20], [1], [6], [3].

From now on, unless otherwise stated, assume that  $U$  is invariant with respect to all isometries in  $E$ , all masses  $m_j > 0$  are positive, and  $U$  is homogeneous of negative degree  $-\alpha$ .

The potential  $U$  is invariant with respect to a suitable subgroup of  $\Sigma_n \times O(E)$ , where  $\Sigma_n$  is the symmetric group on  $n$  elements and  $O(E)$  denotes the orthogonal group on the euclidean space  $E$ . For example, if all masses are equal and  $U$  is defined as in (1.1), then  $G = \Sigma_n \times O(E)$ ; if all masses are distinct, then  $G = \{1\} \times O(E)$ .

The following proposition is a well-known characterization of the set of CC, which we will generalize to relative equilibria in (1.9).

**(1.8)** *Let  $S \subset Y$  denote the inertia ellipsoid, defined as  $S = \{\mathbf{q} \in Y : \|\mathbf{q}\|_M^2 = 1\}$ . A point  $\mathbf{q} \in S$  is a central configuration if and only if it is a critical point of the restriction of  $U$  to  $S$ .*

---

<sup>3</sup>Cf. [18], p. 401.

*Proof.* Critical points of  $U|_S$  are points  $\mathbf{q} \in Y$  such that  $\ker dU \supset T_{\mathbf{q}}S$ . With respect to the (non-degenerate) bilinear form  $\langle -, - \rangle_M$ , this can be written as  $\nabla_M U(\mathbf{q}) = \lambda \nabla_M (\|\mathbf{q}\|_M^2) = 2\lambda \mathbf{q}$ . *q.e.d.*

**(1.9)** Assume  $\dim(E) = 3$ . Let  $C$  be the vertical cylinder defined as

$$C = \{\mathbf{q} \in Y : \langle P\mathbf{q}, \mathbf{q} \rangle_M = c, \}$$

where  $P$  is the projection of  $E$  to  $E'$  as in (1.7) and  $c = \langle P\bar{\mathbf{q}}, \bar{\mathbf{q}} \rangle_M$ . A configuration  $\bar{\mathbf{q}} \in Y$  is a relative equilibrium configuration rotating by  $e^{t\Omega}$  if and only if it is a critical point of  $U$  restricted to  $C \subset X$  and  $U(\bar{\mathbf{q}}) > 0$ .

*Proof.* The configuration  $\mathbf{q}$  is a relative equilibrium configuration if and only if  $\Omega^2 \mathbf{q} = \nabla_M U(\mathbf{q})$ ; since  $\Omega^2 = -\omega^2 P$  for an  $\omega \neq 0$ , this is equivalent to

$$\nabla_M U(\mathbf{q}) = -\omega^2 P\mathbf{q}.$$

On the other hand,  $\nabla_M (\langle P\mathbf{q}, \mathbf{q} \rangle_M) = 2P\mathbf{q}$ , hence  $\mathbf{q} \in C$  is a critical point of the restriction  $U|_C$  iff  $\nabla_M U(\mathbf{q}) = \lambda 2P\mathbf{q}$ . The proof follows since, by homogeneity,

$$\langle \nabla_M U(\mathbf{q}), \mathbf{q} \rangle_M = -\alpha U(\mathbf{q})$$

and  $\langle P\mathbf{q}, \mathbf{q} \rangle_M = \|P\mathbf{q}\|_M^2$ . *q.e.d.*

**(1.10)** Let  $K$  be a subgroup of the symmetry group  $G$  of  $U$  on  $Y$ . Then the inertia ellipsoid  $S$  is  $K$ -invariant, and critical points of the restriction of  $U$  to  $S^K = \{\mathbf{q} \in S : K\mathbf{q} = \mathbf{q}\}$  are precisely the critical points of  $U|_S$  belonging to  $S^K$ . If the vertical cylinder  $C$  is  $K$ -invariant, then critical points of the restriction of  $U$  to  $C^K = \{\mathbf{q} \in C : K\mathbf{q} = \mathbf{q}\}$  are precisely the critical points of  $U|_C$  belonging to  $C^K$ .

*Proof.* It is a consequence of Palais principle of Symmetric Criticality [9]. *q.e.d.*

## 2 Central configurations as (equivariant) fixed points.

In [5] a way to relate planar central configurations to projective classes of fixed points was introduced. We now generalize the results therein to arbitrary dimensions.

Consider a homogeneous potential  $U$ , as above with the further assumption that  $\forall \mathbf{q}, U(\mathbf{q}) > 0$ . From this it follows that  $\nabla_M U(\mathbf{q}) \neq \mathbf{0}$  because  $\langle \nabla_M U(\mathbf{q}), \mathbf{q} \rangle_M = -\alpha U(\mathbf{q})$ .

**(2.1)** The map  $F: S \rightarrow S$  defined as

$$F(\mathbf{q}) = -\frac{\nabla_M U(\mathbf{q})}{\|\nabla_M U(\mathbf{q})\|_M}$$

is well-defined, and a configuration  $\bar{\mathbf{q}} \in S$  is a central configuration if and only if it is a fixed point of  $F$ .

*Proof.* It follows from the assumption that  $\forall \mathbf{q}, \nabla_M U(\mathbf{q}) \neq \mathbf{0}$ , and therefore  $F$  is well-defined. A configuration  $\mathbf{q}$  is fixed by  $F$  if and only if there exists  $\lambda = \|\nabla_M U(\mathbf{q})\|_M > 0$  such that  $\nabla_M U(\mathbf{q}) = -\lambda \mathbf{q}$ . Hence, if  $F(\mathbf{q}) = \mathbf{q}$  then  $\mathbf{q}$  is central. Conversely, by homogeneity of  $U$ ,  $\langle \nabla_M U(\mathbf{q}), \mathbf{q} \rangle_M = -\alpha U(\mathbf{q})$ , and hence if  $\mathbf{q} \in S$  is a central configuration then  $\nabla_M U(\mathbf{q}) = \lambda \mathbf{q} \implies \lambda = -\alpha U(\mathbf{q}) < 0$  and therefore  $F(\mathbf{q}) = \mathbf{q}$ . *q.e.d.*

Now, in general (and without the positivity assumption) the map  $F$  needs not being compactly fixed (see for example Robert's continuum [13] of central configurations with four unit masses and a fifth negative  $-1/4$  mass in the origin<sup>4</sup>). In the gravitational case (positive masses and Newtonian mutual attraction), the map  $F$  turns out to be compactly fixed [5] (see also Shub's estimates [15]).

**(2.2)** *Let  $G$  be the symmetry group of  $U$  on  $Y$ , as above. Then  $F$  is  $G$ -equivariant. For each subgroup  $K \subset G$ ,  $F$  induces a self-map  $\bar{F}: S/K \rightarrow S/K$  on the quotient space  $S/K$ .*

*Proof.* For  $g \in G$ ,  $gS = S$ , and for each  $g$  such that  $U(g\mathbf{q}) = U(\mathbf{q})$  the equality  $\nabla_M U(g\mathbf{q}) = g\nabla_M U(\mathbf{q})$  holds. In fact, since  $U \circ g = U$ ,  $dU \circ g = dU$ , and therefore for each vector  $\mathbf{v}$  one has  $\langle \nabla_M U(\mathbf{q}), \mathbf{v} \rangle_M = dU(\mathbf{q})[\mathbf{v}] = dU(g\mathbf{q})[g\mathbf{v}] = \langle \nabla_M U(g\mathbf{q}), g\mathbf{v} \rangle_M = \langle g^{-1}\nabla_M U(g\mathbf{q}), \mathbf{v} \rangle_M$ . Thus  $F$  is  $G$ -equivariant, and hence  $K$ -equivariant for each subgroup  $K \subset G$ . *q.e.d.*

Let  $U$  be a homogeneous potential with the following property: for each orthogonal projection  $p: E \rightarrow P$  on a plane  $P$ , for each  $\mathbf{q} \in S$  there exists  $j \in \{1, \dots, n\}$  such that

$$(2.3) \quad p\left(\frac{\partial U}{\partial \mathbf{q}_j}(\mathbf{q})\right) \cdot p(\mathbf{q}_j) \leq 0.$$

Moreover, if there exists  $i \in \{1, \dots, n\}$  such that  $p(\mathbf{q}_i) \neq 0$ , then the  $j$  of (2.3) is such that  $p(\mathbf{q}_j) \neq 0$ .

It is easy to see that the Newtonian potential (1.1) (with positive masses and homogeneity  $-\alpha$ ) satisfies Property (2.3): let  $j$  be the index maximizing

---

<sup>4</sup>In the AMS review of [13], D. Saari states that a similar effect occurs for positive masses, but with two or more homogeneous potentials. More precisely, should the potential be a sum of homogeneous potentials, then there always is a continuum of different relative equilibria configurations.

$\|p(\mathbf{q}_k)\|^2$  for  $k = 1, \dots, n$ ; then

$$\begin{aligned}
p \left( \alpha \sum_{k \neq j} \frac{m_j m_k (\mathbf{q}_k - \mathbf{q}_j)}{\|\mathbf{q}_k - \mathbf{q}_j\|^{\alpha+2}} \right) \cdot p(\mathbf{q}_j) &= \alpha \left( \sum_{k \neq j} \frac{m_j m_k (p(\mathbf{q}_k) - p(\mathbf{q}_j))}{\|\mathbf{q}_k - \mathbf{q}_j\|^{\alpha+2}} \right) \cdot p(\mathbf{q}_j) \\
&= \alpha \sum_{k \neq j} \frac{m_j m_k (p(\mathbf{q}_k) \cdot p(\mathbf{q}_j) - \|p(\mathbf{q}_j)\|^2)}{\|\mathbf{q}_k - \mathbf{q}_j\|^{\alpha+2}} \\
&\leq \sum_{k \neq j} \frac{m_j m_k (\|p(\mathbf{q}_k)\| \|p(\mathbf{q}_j)\| - \|p(\mathbf{q}_j)\|^2)}{\|\mathbf{q}_k - \mathbf{q}_j\|^{\alpha+2}} \\
&\leq 0
\end{aligned}$$

It is trivial to see that if  $U$  satisfies (2.3), then the map  $F: S \rightarrow S$  defined in (2.1) satisfies the following property: for each orthogonal projection  $p: E \rightarrow P$  on a plane  $P$ , for each  $\mathbf{q} \in S$ , there exists  $j \in \{1, \dots, n\}$  such that

$$(2.4) \quad p(F_j(\mathbf{q})) \cdot p(\mathbf{q}_j) \geq 0 .$$

**(2.5) Theorem.** *Suppose that  $G \supset SO(E)$ , where as above  $G$  is the symmetry group of  $U$  on  $Y$ , and  $SO(E) = SO(d)$  (the group of rotations in  $E$ ) acting diagonally on  $Y$ . Let  $\overline{F}: S/SO(d) \rightarrow S/SO(d)$  be the map induced on the quotient space by (2.2). Assume that all masses are positive,  $U > 0$  and (2.4) holds. Then, if  $\pi: S \rightarrow S/SO(d)$  denotes the projection on the quotient,*

$$\pi(\text{Fix}(F)) = \text{Fix}(\overline{F}) .$$

*Proof.* If  $\mathbf{q} \in \text{Fix}(F) \subset S$  is fixed by  $F$ , then  $\overline{F}(\pi(\mathbf{q})) = \pi F(\mathbf{q}) = \pi(\mathbf{q}) \implies \pi(\mathbf{q}) \in \text{Fix}(\overline{F})$ , hence  $\pi(\text{Fix}(F)) \subset \text{Fix}(\overline{F})$ . On the other hand, let  $\pi(\mathbf{q}) \in \text{Fix}(\overline{F})$ . Then there exists a rotation  $g \in SO(E)$  such that  $F(\mathbf{q}) = g\mathbf{q}$ . Without loss of generality, we can assume that  $g = e^\Omega$ , where  $\Omega$  is an antisymmetric  $d \times d$  matrix, with  $k$   $(2 \times 2)$ -blocks on the diagonal  $\begin{bmatrix} 0 & -\theta_i \\ \theta_i & 0 \end{bmatrix}$ , with  $\theta_i \in [-\pi, \pi]$  for each  $i = 1, \dots, k$ , and, if  $d$  is odd, a one-dimensional diagonal zero entry ( $d = 2k$  or  $d = 2k + 1$ ). We can also assume that only the first (say,  $l \leq k$ ) blocks have  $\theta_i \neq 0$ , hence  $\Omega$  has  $l$  non-zero  $(2 \times 2)$  diagonal blocks and is zero outside. Note that for each  $\mathbf{x} \in E$ , the quadratic form  $(e^\Omega \mathbf{x}) \cdot (\Omega \mathbf{x})$  on  $E$  can be written with  $l$  non-singular positive-defined blocks

$$\theta_i \begin{bmatrix} \sin \theta_i & -\cos \theta_i \\ \cos \theta_i & \sin \theta_i \end{bmatrix} \sim \theta_i \begin{bmatrix} \sin \theta_i & 0 \\ 0 & \sin \theta_i \end{bmatrix}$$

on the diagonal, and hence it is non-negative. Moreover, if one writes  $\mathbf{x} \in E \cong \mathbb{R}^d$  as  $(\mathbf{z}_1, \mathbf{z}_2, \dots, \mathbf{z}_l, x_{2l+1}, \dots, x_d)$ , with  $\mathbf{z}_i \in \mathbb{R}^2$  for  $i = 1, \dots, l$  and  $x_i \in \mathbb{R}$ , then

$$(2.6) \quad (e^\Omega \mathbf{x}) \cdot (\Omega \mathbf{x}) = \sum_{i=1}^l \theta_i \sin \theta_i \|\mathbf{z}_i\|^2 .$$

Let  $p_i: E \rightarrow \mathbb{R}^2$  denote the projection  $\mathbf{x} \mapsto \mathbf{z}_i$ , for  $i = 1, \dots, l$ .

Since  $F(e^{t\Omega}\mathbf{q})$  does not depend on  $t \in \mathbb{R}$ ,

$$\begin{aligned} 0 &= \frac{d}{dt} \|F(e^{t\Omega}\mathbf{q})\|_M^2|_{t=0} \\ &= 2\langle F(\mathbf{q}), \Omega\mathbf{q} \rangle_M \\ &= 2\langle e^\Omega\mathbf{q}, \Omega\mathbf{q} \rangle_M = \\ &= 2 \sum_{j=1}^n m_j (e^\Omega\mathbf{q}_j) \cdot (\Omega\mathbf{q}_j) \end{aligned}$$

For each  $j = 1, \dots, n$  the inequality  $m_j > 0$  holds, and for each  $\mathbf{x} \in E$  the inequality  $(e^\Omega\mathbf{x}) \cdot (\Omega\mathbf{x}) \geq 0$  holds: it follows that for each  $j$ ,  $(e^\Omega\mathbf{q}_j) \cdot (\Omega\mathbf{q}_j) = 0$ . By (2.6), this implies that, given  $j$ , for each  $i = 1, \dots, l$  either  $p_i(\mathbf{q}_j) = 0$  or  $\theta_i \in \{\pi, -\pi\}$  (since  $\theta_i \neq 0$  for  $i = 1, \dots, l$ ). If  $p_i(\mathbf{q}_j) = 0$  for each  $j$ , then actually  $g\mathbf{q} = \mathbf{q}$ , and therefore  $\pi(\mathbf{q}) \in \pi(\text{Fix}(F))$ . So, without loss of generality one can assume that for each  $i = 1, \dots, l$  there exists  $j$  such that  $p_i(\mathbf{q}_j) \neq 0$ . Suppose that  $l \geq 1$ , and therefore  $\theta_1 = \pm\pi$ . By property (2.4) there exists  $\bar{j}$  such that  $p_1(F_{\bar{j}}(\mathbf{q})) \cdot p_1(\mathbf{q}_{\bar{j}}) \geq 0$  and  $p_1(\mathbf{q}_{\bar{j}}) \neq 0$ . But this would imply that

$$F(\mathbf{q}) = g\mathbf{q} = e^\Omega\mathbf{q} \implies -p_1(\mathbf{q}_{\bar{j}}) \cdot p_1(\mathbf{q}_{\bar{j}}) \geq 0 \implies p_1(\mathbf{q}_{\bar{j}}) = 0,$$

which is not possible. Therefore, if condition (2.4) holds,  $l = 0$ , and  $g\mathbf{q} = \mathbf{q}$ . The conclusion follows. *q. e. d.*

### 3 Projective fixed points and Morse indices

In this section, we finally prove the equation relating fixed point and Morse indices of central configurations.

**(3.1)** *If  $U$  is homogeneous of degree  $-\alpha$ , then for each central configuration  $\mathbf{q}$ , up to a linear change of coordinates*

$$-\alpha U(\mathbf{q})(I - F'(\mathbf{q})) = D^2\tilde{U}(\mathbf{q}),$$

where  $F: S \rightarrow S$  is the function of (2.1), defined as  $F(\mathbf{q}) = -\frac{\nabla_M U(\mathbf{q})}{\|\nabla_M U(\mathbf{q})\|_M}$ , and  $D^2\tilde{U}(\mathbf{q})$  is the Hessian of the restriction  $\tilde{U}$  of  $U$  to  $S$ , evaluated at  $\mathbf{q}$ .

*Proof.* After a linear change of coordinates in  $X$ , we can assume  $m_i = 1$  for each  $i$  and  $\mathbf{q} = (1, 0, \dots, 0) = (1, \mathbf{0})$  (rescale  $\mathbf{q}$  by a diagonal matrix with suitable  $m_j$  on its diagonal, and apply a rotation - this leaves  $U$  homogeneous of the same degree). Given suitable linear coordinates  $\mathbf{x} = (x_0, x_1, \dots, x_k)$  in  $Y \cong \mathbb{R}^{l+1}$ , the ellipsoid  $S$  has equation  $\|\mathbf{x}\|^2 = 1$ , and  $F(\mathbf{x}) = -\frac{\frac{dU}{d\mathbf{x}}}{\|\frac{dU}{d\mathbf{x}}\|}$ ,

where  $\frac{dU}{dx} = \nabla U$  in the  $\mathbf{x}$ -coordinates. Therefore, if  $\mathbf{u} = (u_1, \dots, u_l) \mapsto (\sqrt{1 - \|\mathbf{u}\|}, u_1, \dots, u_l) \in S$  is a local chart around the central configuration  $\mathbf{q} \sim (1, \mathbf{0})$ ,

$$(3.2) \quad \begin{aligned} \frac{\partial F_\alpha}{\partial u_\beta}(\mathbf{0}) &= \frac{\partial^2 U}{\partial x_\alpha \partial x_\beta}(\mathbf{q}) \|\nabla U(\mathbf{q})\|^{-1} \\ D_{\alpha\beta}^2 \tilde{U}(\mathbf{q}) &= \frac{\partial^2 U}{\partial x_\alpha \partial x_\beta}(\mathbf{q}) - \delta_{\alpha\beta} \frac{\partial U}{\partial x_0} . \end{aligned}$$

Now,  $\frac{\partial U}{\partial x_0} = \langle \nabla U, \mathbf{q} \rangle = -\alpha U(\mathbf{q})$ , and since  $\mathbf{q}$  is a central configuration, it is a fixed point of  $F$  and therefore

$$F(\mathbf{q}) = \mathbf{q} \implies \frac{\partial U}{\partial x_\alpha} = 0, \quad \text{for } \alpha = 1, \dots, l$$

and

$$(3.3) \quad \frac{\partial U}{\partial x_0}(\mathbf{q}) = -\|\nabla U(\mathbf{q})\| = -\alpha U(\mathbf{q}) .$$

From (3.3) and (3.2) it follows that for  $\alpha, \beta = 1, \dots, l$ ,

$$(3.4) \quad \begin{aligned} D_{\alpha\beta}^2 \tilde{U}(\mathbf{q}) + \delta_{\alpha\beta} \frac{\partial U}{\partial x_0} &= \frac{\partial^2 U}{\partial x_\alpha \partial x_\beta}(\mathbf{q}) = \frac{\partial F_\alpha}{\partial u_\beta}(\mathbf{0}) \|\nabla U(\mathbf{q})\| \\ \implies D^2 \tilde{U}(\mathbf{q}) &= -\alpha U(\mathbf{q}) I + \alpha U(\mathbf{q}) F'(\mathbf{q}) \end{aligned}$$

*q. e. d.*

**(3.5) Corollary.** *Assume the hypotheses of Theorem (2.5) hold. Then for each non-degenerate projective class of central configurations  $\mathbf{q} \in \text{Fix}(\bar{F})$  in the maximal isotropy stratum of  $\bar{S} = S/SO(d)$ , the fixed point index  $\text{ind}(\mathbf{q}, \bar{F})$  and the Morse index  $\mu_{\tilde{U}}(\mathbf{q})$  of  $\tilde{U}$  at  $\mathbf{q}$  are related by the equality*

$$\text{ind}(\mathbf{q}, \bar{F}) = (-1)^{\mu_{\tilde{U}}(\mathbf{q}) + \epsilon} ,$$

where  $\epsilon = \dim S - d(d-1)/2 + \dim(S) = d(n-1) - 1 - d(d-1)/2$ .

*Proof.* Since  $\tilde{U}$  is  $SO(d)$ -invariant (with diagonal action), and the  $SO(d)$ -orbits in  $S$  with maximal isotropy type have dimension  $d(d-1)/2$ , the Hessian  $D^2 \tilde{U}$  has  $d(d-1)/2$ -dimensional kernel, if  $\mathbf{q}$  is non-degenerate. By proposition (3.1), it follows that  $F'$  has a  $d(d-1)/2$ -multiple eigenvalue 1, and  $\text{ind}(\mathbf{q}, \bar{F}) = (-1)^c$ , where  $c$  is the number of negative eigenvalues of  $-D^2 \tilde{U}$ , i.e. the number of positive eigenvalues of  $D^2 \tilde{U}$ , which is equal to  $\mu_{-\tilde{U}}(\mathbf{q}) = \dim(S - d(d-1)/2) - \mu_{\tilde{U}}(\mathbf{q})$ . *q. e. d.*



**(3.6) Corollary.** *Let  $U$  be the Newton potential in (1.1), with positive masses,  $\alpha > 0$ , and  $\dim(E) = 2$ . Then  $X = \mathbb{C}^n \setminus \Delta$ , and  $S/SO(2) \approx \mathbb{P}^{n-2}(\mathbb{C})_0 \subset \mathbb{P}^{n-2}(\mathbb{C})$ , where  $\mathbb{P}^{n-2}(\mathbb{C})_0$  is the subset of  $\mathbb{P}^{n-1}(\mathbb{C})$  defined in projective coordinates as  $\mathbb{P}^{n-2}(\mathbb{C})_0 \cong \{[z_1 : \dots, z_n] \in \mathbb{P}^n : \sum_j m_j z_j = 0, \forall i, j, z_i \neq z_j\}$ . Then for each non-degenerate projective class of central configurations  $q \in \text{Fix}(\overline{F})$ , with  $\overline{F}: \mathbb{P}^{n-2}(\mathbb{C}) \rightarrow \mathbb{P}^{n-2}(\mathbb{C})$ ,*

$$\text{ind}(q, \overline{F}) = (-1)^{\mu_{\overline{F}}(q)} .$$

*Proof.* If  $d = 2$ , then  $d(n-1) - 1 - d(d-1)/2 = 2(n-2)$ . *q.e.d.*

## 4 An example of non-planar relative equilibrium

**(4.1) Example** (Non-central and non-planar equilibrium solution). In  $E \cong \mathbb{R}^3$ , let  $R_x$ ,  $R_y$  and  $R_z$  denote rotations of angle  $\pi$  around the three coordinate axes. Fix three non-zero constants  $c_1, c_2, c_3$ . Consider the problem with 6 bodies in  $E$ , symmetric with respect to the group  $K$  with non-trivial elements of  $\Sigma_6 \times SO(3)$

$$((34)(56), R_x), ((12)(56), R_y), ((12)(34), R_z) .$$

Assume  $m_j = 1$ , for  $j = 1, \dots, 6$ , and let  $U$  be the potential defined on  $E^6$  by

$$U(\mathbf{q}) = \sum_{i < j} \frac{1 - \gamma_i \gamma_j}{\|\mathbf{q}_i - \mathbf{q}_j\|}$$

where

$$\gamma_1 = \gamma_2 = c_1, \quad \gamma_3 = \gamma_4 = c_2, \quad \gamma_5 = \gamma_6 = c_3 .$$

Now,  $U$  is invariant with respect to  $K$ , and the vertical cylinder is  $K$ -invariant: it follows from (1.9) that critical points of the restriction of  $U$  to  $C^K$  are equilibrium configurations.

In other words, three pairs of bodies of unit masses, each pair charged with charge  $c_j$ , are constrained each pair to belong to one of the coordinate axes and to be symmetric with respect to the origin.

The space  $X^K = Y^K$  has dimension 3, and can be parametrized by  $(x, y, z)$ , where  $x, y$  and  $z$  are (respectively) the coordinates along the corresponding axis of particles 1, 3 and 5. The generic configuration  $\mathbf{q} \in X^K$  can be written as

$$\begin{aligned} \mathbf{q}_1 &= (x, 0, 0) & \mathbf{q}_2 &= (-x, 0, 0) \\ \mathbf{q}_3 &= (0, y, 0) & \mathbf{q}_4 &= (0, -y, 0) \\ \mathbf{q}_5 &= (0, 0, z) & \mathbf{q}_6 &= (0, 0, -z), \end{aligned}$$

and the potential  $U$  restricted to  $X^K$  in such coordinates is

$$U(x, y, z) = \frac{1 - c_1^2}{2|x|} + \frac{1 - c_2^2}{2|y|} + \frac{1 - c_3^2}{2|z|} + 4 \frac{1 - c_1 c_2}{\sqrt{x^2 + y^2}} + 4 \frac{1 - c_1 c_3}{\sqrt{x^2 + z^2}} + 4 \frac{1 - c_2 c_3}{\sqrt{y^2 + z^2}} .$$

The vertical cylinder  $C^K \subset X^K$  is

$$\begin{aligned} C^K &= \{\mathbf{q} \in X^K : \langle P\mathbf{q}, \mathbf{q} \rangle_M = 2\} \\ &= \{(x, y, z) \in X^K : 2x^2 + 2y^2 = 2\} . \end{aligned}$$

Hence an equilibrium solution is a critical point of  $U(\cos t, \sin t, z)$  with positive value  $U > 0$ . Now, assume

$$(4.2) \quad c_1 > 1, \quad c_2 < -1, \quad c_3 < -1 .$$

Then  $1 - c_1^2 < 0$ ,  $1 - c_2^2 < 0$ ,  $1 - c_3^2 < 0$ ,  $1 - c_2c_3 < 0$ ; moreover,  $1 - c_1c_2 > 0$  and  $1 - c_1c_3 > 0$ . The restricted potential  $U$  is defined as

$$U(t, z) = \frac{1 - c_1^2}{2|\cos t|} + \frac{1 - c_2^2}{2|\sin t|} + \frac{1 - c_3^2}{2z} + 4(1 - c_1c_2) + 4\frac{1 - c_1c_3}{\sqrt{\cos^2 t + z^2}} + 4\frac{1 - c_2c_3}{\sqrt{\sin^2 t + z^2}},$$

and is defined on the strip  $(t, z) \in T = (0, \pi/2) \times (0, +\infty)$ .

If

$$(4.3) \quad 1 - c_1^2 + 8(1 - c_1c_3) < 0,$$

then for each  $(t, z)$

$$\begin{aligned} \frac{1 - c_1^2}{2|\cos t|} + 4\frac{1 - c_1c_3}{\sqrt{\cos^2 t + z^2}} &< \frac{1 - c_1^2}{2|\cos t|} + 4\frac{1 - c_1c_3}{|\cos t|} \\ &= \frac{1}{2|\cos t|} (1 - c_1^2 + 8(1 - c_1c_3)) < 0 \end{aligned}$$

and hence  $U(t, z) < 4(1 - c_1c_2)$  for each  $(t, z) \in T$  and  $U \rightarrow -\infty$  on the boundary of  $T$ .

Furthermore, for each  $t \in (0, \pi/2)$ ,

$$\begin{aligned} \frac{\partial U}{\partial z} &= \frac{c_3^2 - 1}{2z^2} + 4z\frac{c_2c_3 - 1}{(z^2 + \sin^2 t)^{3/2}} - 4z\frac{1 - c_1c_3}{(z^2 + \cos^2 t)^{3/2}} \\ &< \frac{c_3^2 - 1}{2z^2} + 4\frac{c_2c_3 - 1}{z^2} - 4z\frac{1 - c_1c_3}{(z^2 + 1)^{3/2}} . \end{aligned}$$

It follows that if

$$(4.4) \quad c_3^2 - 1 + 8(c_2c_3 - 1) < 8(1 - c_1c_3),$$

then for every  $t \in (0, \pi/2)$  there exists  $z_0$  such that  $z > z_0 \implies \frac{\partial U}{\partial z}(t, z) < 0$ .

Thus, whenever both (4.3) and (4.4) hold, the restriction  $U(t, z)$  attains its maximum in the interior of the strip  $T$ . Such a maximum  $(t_m, z_m)$  corresponds to an equilibrium configuration if and only if  $U(t_m, z_m) > 0$ , from (1.9). Note that if  $t = \pi/4$  and  $z = 1/2$  then

$$U(\pi/4, z) = \frac{\sqrt{2}}{2}(2 - c_1^2 - c_2^2) + 4(1 - c_1c_2) + 1 - c_3^2 + \frac{4(2 - c_3(c_1 + c_2))}{\sqrt{3/4}},$$

which is positive, for example, if  $c_1 = 20$  and  $c_2 = c_3 = -2$ . Such coefficients satisfy (4.2), (4.3) and (4.4), and therefore for such a choice of  $c_i$  there exist non-planar relative equilibrium configurations.

## References

- [1] ALBOUY, A. Symétrie des configurations centrales de quatre corps. *C. R. Acad. Sci. Paris Sér. I Math.* 320, 2 (1995), 217–220.
- [2] ALBOUY, A., AND CHENCINER, A. Le problème des  $n$  corps et les distances mutuelles. *Invent. Math.* 131, 1 (1998), 151–184.
- [3] ALBOUY, A., AND KALOSHIN, V. Finiteness of central configurations of five bodies in the plane. *Ann. of Math. (2)* 176, 1 (2012), 535–588.
- [4] BETTI, E. Sopra il moto di un sistema di un numero qualunque di punti che si attraggono o si respingono tra loro. *Annali di Matematica Pura ed Applicata (1867-1897)* 8, 1 (1877), 301–311.
- [5] FERRARIO, D. L. Planar central configurations as fixed points. *J. Fixed Point Theory Appl.* 2, 2 (2007), 277–291.
- [6] HAMPTON, M., AND MOECKEL, R. Finiteness of relative equilibria of the four-body problem. *Invent. Math.* 163, 2 (2006), 289–312.
- [7] MOECKEL, R. On central configurations. *Math. Z.* 205, 4 (1990), 499–517.
- [8] PACELLA, F. Central configurations of the  $N$ -body problem via equivariant Morse theory. *Arch. Rational Mech. Anal.* 97, 1 (1987), 59–74.
- [9] PALAIS, R. S. The principle of symmetric criticality. *Comm. Math. Phys.* 69, 1 (1979), 19–30.
- [10] PALMORE, J. I. Classifying relative equilibria. I. *Bull. Amer. Math. Soc.* 79 (1973), 904–908.
- [11] PALMORE, J. I. Classifying relative equilibria. II. *Bull. Amer. Math. Soc.* 81 (1975), 489–491.
- [12] PALMORE, J. I. Classifying relative equilibria. III. *Lett. Math. Phys.* 1, 1 (1975/76), 71–73.
- [13] ROBERTS, G. E. A continuum of relative equilibria in the five-body problem. *Phys. D* 127, 3-4 (1999), 141–145.
- [14] SAARI, D. G. *Collisions, rings, and other Newtonian  $N$ -body problems*, vol. 104 of *CBMS Regional Conference Series in Mathematics*. Published for the Conference Board of the Mathematical Sciences, Washington, DC; by the American Mathematical Society, Providence, RI, 2005.

- [15] SHUB, M. Appendix to Smale's paper: "Diagonals and relative equilibria". In *Manifolds – Amsterdam 1970 (Proc. Nuffic Summer School)*, Lecture Notes in Mathematics, Vol. 197. Springer, Berlin, 1971, pp. 199–201.
- [16] SMALE, S. Topology and mechanics. I. *Invent. Math.* 10 (1970), 305–331.
- [17] SMALE, S. Topology and mechanics. II. The planar  $n$ -body problem. *Invent. Math.* 11 (1970), 45–64.
- [18] SPIVAK, M. *A comprehensive introduction to differential geometry. Vol. I*, second ed. Publish or Perish, Inc., Wilmington, Del., 1979.
- [19] WINTNER, A. *The Analytical Foundations of Celestial Mechanics*. Princeton Mathematical Series, v. 5. Princeton University Press, Princeton, N. J., 1941.
- [20] XIA, Z. Central configurations with many small masses. *J. Differential Equations* 91, 1 (1991), 168–179.