

Mr Chairman, Ladies & gentlemen!

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The distributed parameter system I am considering is DPS formed by a parabolic PDE, defined in a space-time domain.  $x$  represents the scalar space variable.

Already at this point many questions arise:

- 1) Why distributed parameters i.e. infinite dimensional systems at all?

If time is left, I leave a tentative answer to this question at the end of the talk.

Let me proceed. I want to identify the leading coefficient (conductivity) from knowledge of the (potential, source term) data pair.

The next question is: why inverse problems for PDES?

Well, for the same reasons parameter identification arises with finite dimensional systems.

Almost all inverse problems are known to be ill-posed (otherwise they would not be so interesting). Existence, uniqueness and stability have to be obtained at a price.

I have mentioned stability! In this context, stability will denote a property of the map which relates the data to the parameters. Stability here is a synonym for the type of continuity exhibited by the map.

DEF

In detail, consider the space-time domain  $\Omega$  and the parab. eqn., which shall be understood in the distribution sense.

The inverse problem consists of finding the conductivity  $a$  in the admissible set  $A$ , given  $f$  in this space and  $u \in X$ . Note that  $X$  is made of continuous trajectories s.t. spatial derivatives are in  $L^2$  up to the 2nd order.

The admissible set  $A_{ad}$  is made of bounded, measurable functions, which are bounded away from 0 and from above by given constants.

The key technical tool in the following is the defect equation, also to be understood in the distribution sense, which relates the difference between the 2nd & the reference potential to the difference  $B$  between the 2nd and the ref conductivity.

It is assumed throughout that admissible conductivities  $a, b$  exist. This is important.

The purpose of this talk is to provide a unifying UNI view on both uniqueness and stability of the solution to the considered direct problem.

Uniqueness is obtained by providing a Cauchy datum for the defect equation, which is an ODE wrt to  $B$ .

In general Cauchy data may be available either at regular or singular points.

A singular point at where the partial derivative of potential wrt  $x$  vanishes.

### TREG

I can now state the uniqueness result for the full Cauchy problem - introduce the defect  $\tau$  and this quotient, which makes sure, provided the denominator does not vanish. I define  $B_{ad}$  by requiring that  $B$  be continuous at the left endpt of the interval and vanish there. The uniqueness result perhaps is more easily read when written:

It requires some conditions to be met. By the way,  
 $R^{t-1}$  is the <sup>spatial</sup> antiderivative of a distribution, since all  
antider. of a distribution differ by a constant, then  
antidifferentiation is a set valued map, as it is for  
ordinary functions - Now, these antiderivatives shall be  
continuous @  $x_0^+$ .

If at the instant  $t$  these hypotheses hold, then  $B$  is  
zero a.e. in  $D$ .

USIN.

The uniqueness statement corresponding to the singular problem  
relies on substantially different hypotheses - more precisely it  
is based on the properties of the set of pts where potential is  
spatially stationary.

I assume that in addition to  $\tilde{A}$  (the reference and)  
there are another admissible and  $\tilde{A}'$  s.t. the corresponding  
potentials coincide a.e. in  $\Omega$ .

Here  $f$  can be less regular than specified before in fact, also  
 $u$  may be less regular.  $V_{\tilde{A}'}d$  is a space larger than  $X$ .  
I do not give the details, here and refer you to the proceedings.  
Uniqueness is obtained

either if at a given  $t$ ,  $E_u(t)$  is non-empty but has  
zero Lebesgue measure

or if  $E_u$  is never empty and the measure of  
this intersection is zero.

Here is another condition, which is non-local and which I call self-identifiability - It may be  
interesting to examine it, however, I skip it for the  
sake of conciseness.

When it comes to stability, additional constraints have to be introduced i.e., the problem must be  
regularized.

Assume that uniqueness comes from a regular Cauchy problem.  
The previous H<sub>p</sub> are not sufficient - more work must be  
added - these reciprocals shall e.g. be bounded above by  
given constants - this is an example of regularization.  
Similarly for this L<sup>2</sup> norm.

But the most relaxed restriction affects the reference conductiv-  
ity. It shall be piecewise regular, although it  
may have a countable infinity of discontinuity pts.

Here  $\delta(\cdot)$  is Sidai's measure. Moreover, the series  
of these abs values shall converge  
under these H<sub>p</sub>. There is a uniform (i.e. L<sup>0</sup>) stability  
estimate for B'. The L<sup>0</sup> norm of B is bounded above  
by the product of some constants times the X<sub>T</sub>-norm  
of V.

In fact, the quantity between square brackets can be replaced  
by a suitable constant - what matters is that it only de-  
pends on the reference conductive,  $\tilde{\alpha}$ , not on  $\alpha$ .

The H<sub>p</sub> which affect conductivity may seem restrictive APPL-  
wise. However, there are many applications where they are  
met. Consider e.g. a layered medium. Conductivity  
satisfying this condition is such that the anti derivatives  
of the abs value of  $\alpha'$  are all in L<sup>0</sup>. Moreover, the reference  
 $\tilde{\alpha}$  is continuous at  $x_0$  and  $x_1$ .

### SSIN

If uniqueness is due to a singular Cauchy problem,  
we shall expect two things

- i) that an estimate be obtained, provided the reference  
conductivity is continuous at the singular point
- ii) that the type of stab estimate differ from the  
previous one, in spite of increased local regularity

of  $\hat{a}$ ,

I have to introduce the closed interval  $\bar{I}$  - I require that  $\forall u_x$  be in  $L^1$ , at least at that instant of time  $t$ , when the set  $E_u$  complies with this condition, which is not sufficient to yield uniqueness.

Now the interesting fact is that  $E_v(t)$  may have positive measure. Both sets shall be contained in  $\bar{I}$ , however

$\forall v_x$  shall be in  $L^1$  of this interval - I assume  $\hat{a}$ , besides being admissible, is continuous on  $\bar{I}$  and that  $b$ , besides being admissible, is made to coincide with  $\hat{a}$  inside  $E_v(t)$ . Namely  $b$  would not be uniquely determined there.

Under these circumstances, there is an  $L^1$  estimate for  $B$ , where the estimation constant is the product of  $C_v$ , introduced here, and the quantity  $1 +$  twice the  $L^\infty$  norm of  $\hat{a}$  - this norm can be replaced by  $a_H$ .

## SYN

Why do stability inequalities differ?

In both cases the starting pt is the defect equation at time  $t$ .

The regular Cauchy pt corresponds to this condition:  $B(x_0)$  vanishes. I can then divide all terms in the eqn by  $v_s$ , which vanishes nowhere (as required by the H<sub>p</sub>), then estimate the growth of  $B$  - standard stability theory for ODEs, based on a generalization of Gronwall-Bellman's inequality, yields this estimate, where the antiderivative of  $R$  must be rewritten as a sum of  $V$  - this eventually yields the  $L^\infty$  estimate you have seen a few slides ago.

on the other hand, if uniqueness comes from a unique Cauchy problem, the quantity which vanishes at time  $T$  and at the point  $S(T)$  is this product.

I cannot divide by  $v_x(\cdot)$  and carry out the same estimate as in the previous case, because  $v_x$  is uniformly zero at some point.

I do divide by  $v_x$ , but then I have to integrate the quotient over the interval  $D$ , which gives me an  $L^1$ -estimate for  $B$ .

The remaining step is to replace this antiderivative by a function of  $V$ .

We have seen the role played by the defect CONC equation and the related Cauchy problems. They lead to a comprehensive view of these uniqueness & stability estimates (conditions).

Also, I have shown the role of regularization conditions required to attain stability estimates.

I hope the synopsis of the proof has given an idea of the technical details.

The last remark in the list is about the finite independent case: everything applies in a straightforward way.

I can now answer the question: why infinite dimensional systems? In practical simulations all PDEs are replaced by algebraic systems. It may seem that finite dim systems suffice. On the other hand we often want to refine the grid over which our algebraic system is defined. We want therefore to predict what happens as the grid is made finer & finer. This implies passing to the limit, hence investigating an infinite-dimensional system, governed by a PDE. Thank you!